

# Study of AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMTs with TiO<sub>2</sub> gate dielectric and regrown source/drain

Z. Touati<sup>a</sup>, Z. Hamaizia<sup>a</sup> and Z. Messai<sup>b,c</sup>

<sup>a</sup>Laboratory of Semiconducting and Metallic Materials, University of Mohamed Khider Biskra, Algeria.

<sup>b</sup>Electronics Department, Faculty of Sciences and Technology, University of BBA, Algeria

<sup>c</sup>Laboratory of Optoelectronics and Components, UFAS 19000, Algeria

\*Corresponding author, email: zinouu113@yahoo.fr

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## Abstract

In this work, we proposed a novel TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> metal-oxide-semiconductor high electron mobility transistor model (MOS-HEMT) with 60 nm gate-length and high-k TiO<sub>2</sub> gate dielectric. The DC and RF characteristics of the proposed AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT structure were obtained using the TCAD Silvaco Software. It shows a maximum extrinsic transconductance of 198.3 mS/mm, a saturated drain-current density at V<sub>gs</sub> = 4 V of 668.4 mA/mm, with a maximum of 677.9 mA/mm, a unity-gain cut-off frequency of 229.8 GHz, and an impressive maximum oscillation frequency of 627.8 GHz. Characterized at 10 GHz, the power performance of the proposed device shows an output power of 22.3 dBm, a power gain of 13.1 dB, and a power-added efficiency of 26.5%. The obtained simulated results are very encouraging while compared to existing AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMTs. This is, in fact, the best TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT simulated high-frequency performance reported so far, making it suitable for high-power RF circuit applications.

**Keywords:** TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub>; MOS-HEMT; high-k, TiO<sub>2</sub>; Regrown Source/Drain; SILVACO.

## 1. Introduction

GaN-based High Electron Mobility Transistors (HEMTs) are quite suitable devices for high-power and high-frequency applications, due to their appropriate material properties such as high breakdown voltage, high saturation velocity, low effective mass, high thermal conductivity and high two-dimensional electron gas (2DEG) density, in the order of 10<sup>23</sup> cm<sup>-2</sup>, produced by spontaneous and piezoelectric polarization in AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures [1-3]. However, Schottky gate transistors usually exhibit high gate leakage current [4], making the drain current to collapse when operating at high frequencies [5] thus, limiting the performance and reliability of HEMTs in Radio Frequency (RF) power applications.

Metal Oxide Semiconductor HEMTs (MOS-HEMTs) with insulating dielectric have been widely investigated, and excellent performance have been reported using various high-k insulators (Al<sub>2</sub>O<sub>3</sub> [4-6], TiO<sub>2</sub> [7-9], HfO<sub>2</sub> [10, 11], Pr<sub>2</sub>O<sub>3</sub> [12, 13], SiN [14], SiO<sub>2</sub> [14] and NiO [15]).

In [16], Hong-Yu et al. asserted that SiO<sub>2</sub>-gate AlGa<sub>N</sub>/Ga<sub>N</sub> HEMTs show good electrical performance such as high unity-gain cut-off frequency f<sub>u</sub> and maximum oscillation frequency f<sub>max</sub>, low gate leakage, small threshold voltage V<sub>th</sub> fluctuation and small current collapse. However, high leakage currents were observed in MOS-HEMTs under forward bias conditions [17-19], probably due

to insufficient AlGa<sub>N</sub> bandgaps [16]. In addition, Deen et al. [17] and Hayashi et al. [20] reported severe V<sub>th</sub> fluctuations from HEMTs with HfO<sub>2</sub>-gate dielectric materials while Stoklas et al. [21] described similar instability issues in HEMTs with ZrO<sub>2</sub>-gate dielectric materials. Therefore, further investigation is necessary to gain better understanding of high-k insulators/(Al)Ga<sub>N</sub> interfaces.

Furthermore, all these devices suffer from a high contact resistance (of about 0.3 Ω·mm) and high on-resistance (of about 1 Ω·mm) due to the alloyed ohmic contacts and large source-drain distance. Recently, n<sup>+</sup> regrown source/drain regions have become a novel approach leading to non-alloyed ohmic contacts and a self-aligned gate process [22, 23].

In this work, the DC and microwave characteristics of a novel AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT with 5-nm-thick TiO<sub>2</sub> dielectric layer, 60 nm gate-length, and 10 nm heavily doped (n<sup>++</sup> InGa<sub>N</sub>) regrown ohmic source/drain regions was investigated.

Using the commercial TCAD Silvaco Software, the proposed TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT device exhibits a maximum extrinsic transconductance (g<sub>m</sub>) of 198.3 mS/mm, a saturated I<sub>ds</sub> density of 668.4 mA/mm at V<sub>gs</sub> = 4 V, with a maximum at 677.9 mA/mm, a unity-gain cut-off frequency (f<sub>u</sub>) of 229.8 GHz, as well as an impressive maximum oscillation frequency f<sub>max</sub> of 627.8 GHz. In terms of power performance, it can deliver an output power (P<sub>out</sub>) of 22.3

dBm, a power gain of 13.1 dB, and a power-added efficiency (PAE) of 26.5% at 10 GHz, making it suitable for high-power RF circuit applications.

## 2. Device description and simulation models

### 2.1. The oxide choice

As reported in Table.1, among various gate dielectric materials[24], TiO<sub>2</sub> can be considered as the most suitable candidate because of its large static dielectric constant (k= 86–170). Previous research works[25-27] showed that transistors, with TiO<sub>2</sub> as gate dielectric, demonstrate a high breakdown voltage and very low gate leakage current, along with a slight decrease in  $g_m$  and small shift in  $V_{th}$ .

Table 1: High-k dielectric materials and their properties[24].

Gate dielectric Material	Dielectric constant (k)	Energy bandgap Eg (eV)	Conduction band offset $\Delta E_c$ (eV)	Valence band offset $\Delta E_v$ (eV)
SiO <sub>2</sub>	3.9	9	3.5	4.4
Al <sub>2</sub> O <sub>3</sub>	8	8.8	3	4.7
TiO <sub>2</sub>	80	3.5	1.1	1.3
ZrO <sub>2</sub>	25	5.8	1.4	3.3
HfO <sub>2</sub>	25	5.8	1.4	3.3

### 2.2. The device structure

Figure.1 shows the cross-sectional schematic of the proposed TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/GaN MOS-HEMT device with source/drain regrowth. The gate length and width are of 60 nm and 20 nm respectively, the source/drain length is of 0.5  $\mu$ m, while both source-gate and gate drain spacing are of 0.33  $\mu$ m. The oxide thickness is 5 nm with TiO<sub>2</sub> dielectric to minimize leakages. A GaN buffer layer of 0.8  $\mu$ m thickness (doped with [donor] =  $2.5 \times 10^{18}$  cm<sup>-3</sup>) was adopted and grown on top of the SiC semi-insulating substrate.

An AlN nucleation layer has been inserted to reduce stress and lattice mismatch, with an undoped GaN channel thickness of 20 nm. A 7 nm Al<sub>0.3</sub>GaN<sub>0.7</sub> barrier layer, which depletes the 2DEG, provides a strong carrier confinement in the quantum well at the hetero-interface and minimizes the junction leakage and the off-state leakage current  $I_{off}$ .

Next, a graded n+ In<sub>0.3</sub>GaN<sub>0.7</sub>N (10 nm), doped with [donor] =  $2 \times 10^{19}$  cm<sup>-3</sup>, was introduced to reduce the access and contact resistances[22]. Non-alloyed contacts were formed for the source/drain regions to give a low contact resistance. A passivation of the SiO<sub>2</sub> layer over the dielectric surface

was used to reduce the surface trap density to further improve the device characteristics.

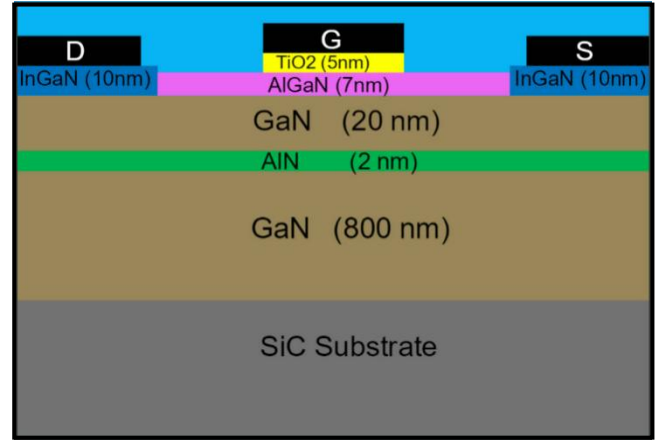


Figure 1. Cross-section structure of the proposed TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/GaN MOS-HEMT.

### 2.3. Physical models

Simulations were performed using the Two-Dimensional (2D) Silvaco ATLAS TCAD tool. The Boltzmann transport theory has shown that the current densities in the continuity equations may be approximated by a drift-diffusion model (DD). This model is one of the most basic carrier transport models in semiconductor physics. In this case, the current densities for electrons and holes under the DD model can be expressed as

$$\vec{J}_n = -nq\mu_n \nabla \Phi_n \quad (1)$$

$$\vec{J}_p = -nq\mu_p \nabla \Phi_p \quad (2)$$

where n and p state for the respective electron and hole concentrations, and  $\mu_n$  and  $\mu_p$  for the respective electron and hole mobilities.  $\Phi_n$  and  $\Phi_p$  represent the respective electron and hole quasi-fermi potentials.

Based on the DD model, the Poisson equation (3), the electron continuity equation (4) and the hole continuity equation (5), were numerically solved[28]. A drift-diffusion model was used to solve the transport equation.

$$\text{div}(\epsilon \nabla \Psi) = -\rho \quad (3)$$

$$\frac{dn}{dx} = \frac{1}{q} \nabla \vec{J}_n + G_n - R_n \quad (4)$$

$$\frac{dp}{dx} = \frac{1}{q} \nabla \vec{J}_p + G_p - R_p \quad (5)$$

with the permittivity,  $\Psi$  the electrostatic potential, q the electron charge, and  $\rho$  the space charge density. ( $\vec{J}_n, G_n, R_n$ )

and ( $\vec{J}_p$ ,  $G_p$ ,  $R_p$ )state for the current densities, generation rates, and recombination rates for respectively the electrons and holes.

The considered physical models in the simulations include Shockley-Read Hall recombination (SRH), Fermi-Dirac statistics, Multiband k.p., GaNFET model and field-dependent mobility [29].

### 3. Results

#### 3.1. Energy band diagram of MOS-HEMT

Figure.2 shows the energy band of the TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/Ga<sub>N</sub> MOS-HEMT biased at  $V_{GS} = 1$  V and  $V_{DS} = 0.5$  V. When  $V_{GS}$  is swept from -10 V to 1 V, the captured electrons begin to be thermally emitted into the AlGa<sub>N</sub> conduction band. The blue arrows show a possible path of the thermally emitted electrons from the capture states: (1) electrons emitted from trap states to the conduction band of the AlGa<sub>N</sub> layer, at the TiO<sub>2</sub> /AlGa<sub>N</sub> interface; (2) electrons at the AlGa<sub>N</sub> conduction band 'rolling down' to the triangular quantum well at the AlGa<sub>N</sub>/Ga<sub>N</sub> interface due to the potential difference across the thickness of the AlGa<sub>N</sub> layer. This causes the 2DEG to appear at the AlGa<sub>N</sub>/Ga<sub>N</sub> interface; (3) electrons are swept away laterally due to  $V_{DS}$ .

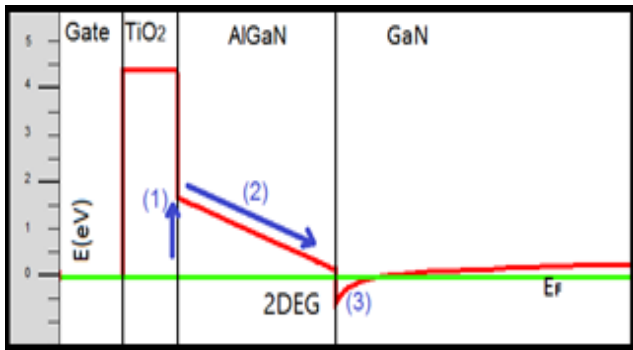


Figure 2. Energy band of the TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/Ga<sub>N</sub> MOS-HEMT under the gate electrode.

#### 3.2. DC results

DC simulations were performed in Silvaco by sweeping the gate voltage from -10 V to 10 V and the drain voltage from 0 V to 6 V. The  $I_{DS}$ - $V_{DS}$  curves of Figure.3 allowed evaluating some MOS-HEMT parameters such as the knee voltage (transition between the triode and saturation regions), the on-resistance, as well as the maximum current and self-heating. The device exhibited a peak current density of 638 mA/mm at  $V_{GS} = 2$  V and 687 mA/mm at  $V_{GS}=10$  V. However, the device did not pinch off completely even at a gate bias of -6V. As documented by Breitschadel et al [30],

Zhongda et al [31] and Park et al [32], this is probably due to the short-channel effect, since the simulated device has a gate length of 60 nm.

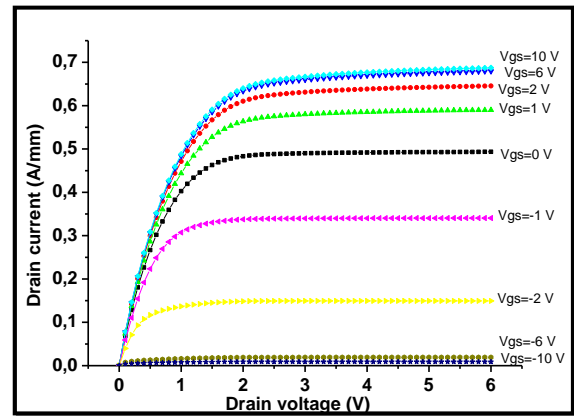


Figure 3.  $I_{DS}$ - $V_{DS}$  DC characteristics of the simulated TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/Ga<sub>N</sub> MOS-HEMT.

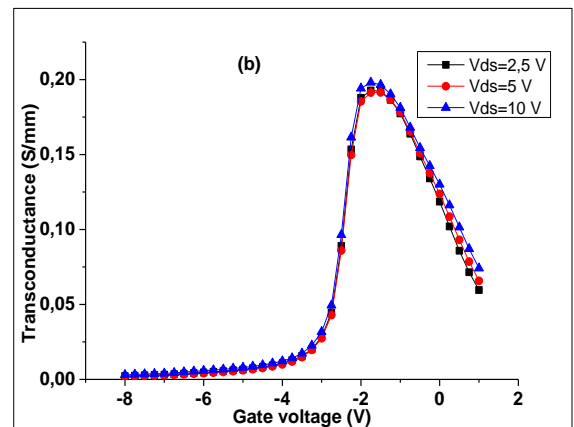
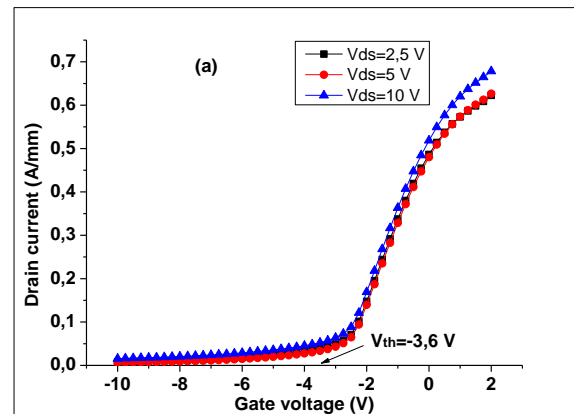


Figure 4. (a) Transfer characteristic, (b) transconductance at  $V_{DS} = 2.5$  V, 5 V and 10 V.

As shown in Figure 4 (a), the threshold voltage  $V_{th}$  is about -3.6 V. The transconductance  $g_m$  (Figure.4 (b)) was obtained by deriving the  $I_{ds}$ - $V_{cs}$  curves at fixed  $V_{ds}$ , showing a peak of about 200 mS/mm.

3.3. Analysis of traps and defects

Traps can be caused by defects in the material, that can be related to crystal defects, polluting agents and mismatch between different layers[33]. The trapping process results also in strong dependencies of the I-V characteristics in drain-source voltage slowdynamics [34].

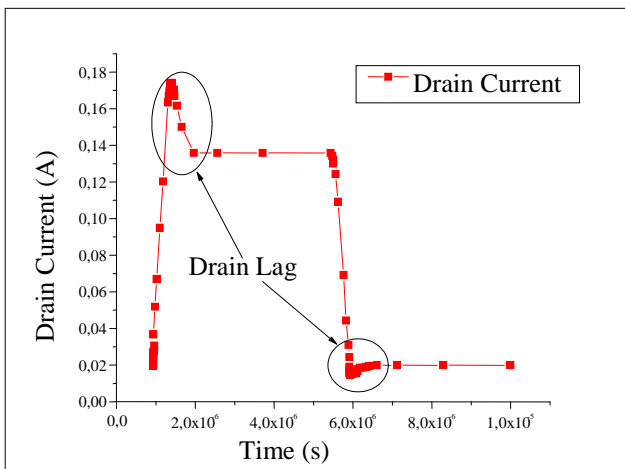
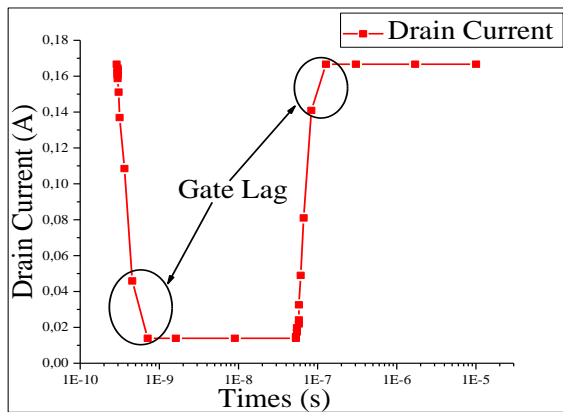


Figure 5. (a) Gate Lag and (b) Drain Lag of the TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT.

From an electrical point of view, when characterizing TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/Ga<sub>N</sub> MOS-HEMT, the traps can be separated into two types “Gate Lag” and “Drain Lag” [34], corresponding to the delayed response of the drain current associated with the  $V_{cs}$  and  $V_{ds}$  control voltages, respectively.

Note that drain-lag effects can be considered as self-back gating effects caused by deep traps under the channel.

Since each effect depends only on its respective drain or gate voltage, in frequency domain, the impact of gate-lag is noticeable on the transconductance (the partial derivative of the drain current with respect to  $V_{cs}$ ) while the impact of drain-lag is focused on the output conductance (the partial derivative of the drain current with respect to  $V_{ds}$ ) [35]. Analysis of the drain- and gate-lag effects is a good indicator of the device RF performance. A device with reduced gate- and drain-lag can be accurately predicted to show improved RF performance, with higher  $f_{rf}$  and  $f_{max}$ . According to the simulation results shown in Figure.5, it is clear that the Ga<sub>N</sub> MOS-HEMT examined in this work showed both gate- and drain-lag.

The calculated gate-lag and drain-lag percentages are given in Table.2. Since our device displays an acceptable drain-lag of 13.6 % and gate-lag of 5.3 %, we can then conclude that our device reaches its steady state promptly.

Table 2: Calculated gate-lag and drain-lag percentage.

Gate-lag (%)	Drain-lag (%)
5.3	13.6

3.4. Impact ionization:

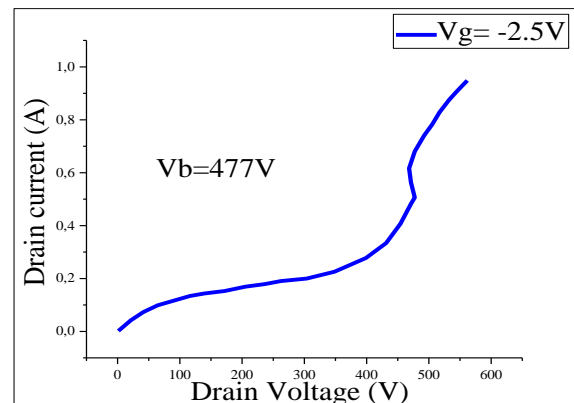


Figure 6. Off-state breakdown voltage simulation results.

The cause of the breakdown is defined as avalanche generation (impact ionization). The breakdown is due to the charge multiplication that occurs when the width of a space charge region is greater than the mean free path between two ionizing impacts. The reciprocal of the mean free path is called the ionization coefficient [36].

In this section, the gate voltage was kept at -2.5 V in order to have a very small drain current and simulate the off-state of the MOS-HEMT. The drain voltage was increased until the break criteria was reached. We can see from Figure.6 that for small values of the drain voltage, the increase in drain current is negligible. However, beyond around 450V, even small variations of the drain voltage induce a large increase in the drain current. This is due to the breakdown of the device. Therefore, we can conclude that the breakdown voltage of the simulated model is close to 477V.

### 3.5. High-frequency results

The frequency performance of microwave devices can be evaluated through small-signal S-parameter simulations, from which several useful pieces of information can be extracted.

Cut off frequency  $f_r$  and maximum oscillation frequency  $f_{max}$  are two of them, representing important figures of merit as far as the frequency limits of the device is concerned.  $f_r$  is defined as the frequency at which forward current gain ( $H_{21}$ ) from hybrid parameters becomes unity and  $f_{max}$  is defined as the frequency where the unilateral gain ( $U_g$ ) or maximum stable gain (MSG) becomes unity [4]. The gains  $H_{21}$ ,  $U_g$  and MSG were extracted directly from the simulated S-parameter data using the following equations.

$$H_{21} = \left| \frac{-2 S_{21}}{(1-S_{11}) \times (2-S_{22}) - S_{12} \times S_{21}} \right| \tag{4}$$

$$U_g = \frac{|S_{21}|^2}{(1-|S_{11}|^2) \times (1-|S_{22}|^2)} \tag{5}$$

$$MSG = \frac{|S_{21}|}{|S_{12}|} \times (K \pm \sqrt{K^2 - 1}) \tag{6}$$

with

$$K = \frac{1-|S_{11}|^2-|S_{22}|^2+|S_{11}S_{22}-S_{12}S_{21}|^2}{2 \times |S_{12}|^2 \times |S_{22}|^2} \tag{7}$$

the stability factor.

Figure.7 shows the simulated S-parameters of the device at bias point  $V_{gs} = -2$  V and  $V_{ds} = 15$  V. The Simulation frequency range was from 0.1 GHz to 20.1 GHz, step-size of 0.1 GHz. Using the S-parameter simulation data and equations (4),(5), (6), and (7), the current gain  $H_{21}$ , the unilateral gain  $U_g$  and the Maximum Stable Gain (MSG) could be calculated as a function of frequency. Figure.8 shows the small signal characteristics of the proposed TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/GaIn MOS-HEMT device.

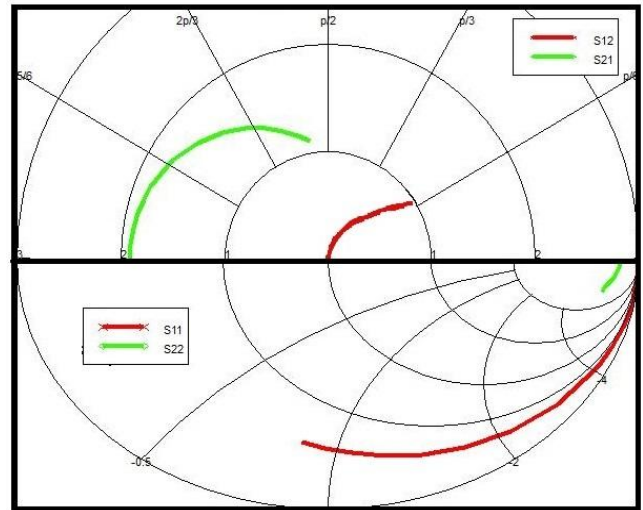


Figure 7. Typical simulated S-parameters of the MOS-HEMT.

The extraction process of  $f_r$  and  $f_{max}$  involves the S-parameters of the transistor at a selected frequency range (here from 0.1 GHz to 20.1 GHz) and a given bias point ( $V_{gs} = -2$  V and  $V_{ds} = 15$  V). Knowing that  $f_r$  is the frequency value where  $H_{21}$  becomes 0 dB and  $f_{max}$  where  $U_g$  or MSG becomes 0 dB [37], we found that  $f_r = 229.83$  GHz and  $f_{max} = 627.58$  GHz, an impressive value.

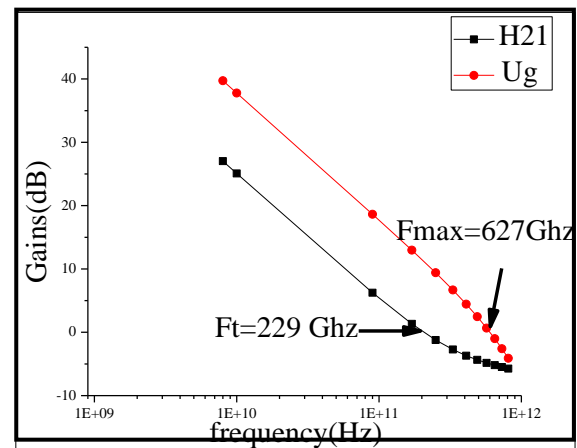


Figure 8. Small signal characteristics of the TiO<sub>2</sub>/Al<sub>0.3</sub>GaN<sub>0.7</sub>/GaIn MOS-HEMT at the bias point  $V_{gs} = -2$  V and  $V_{ds} = 15$  V.

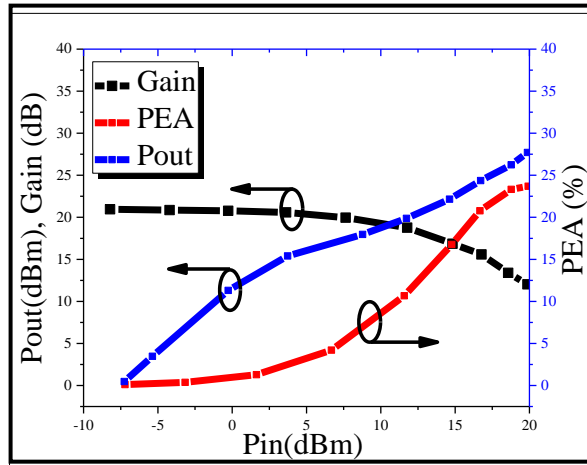


Figure 9. Power characteristics, ( $P_{out}$ , Gain and PAE) of the  $TiO_2/AlGaIn/GaN$  MOS-HEMT at 10 GHz.

To characterize the power performance of the  $TiO_2/AlGaIn/GaN$  MOS-HEMT at 10 GHz, ten large signal input amplitudes were defined using WAVEFORMS statements. Each of these WAVEFORMS were applied to the gate in order to increase the amplitude by using the LOOP statements [29]. Figure.9 shows the simulated typical output power and Power Added Efficiency (PAE) of the device. Table.3 reports the power characteristics of the simulated  $TiO_2/Al_{0.3}Ga_{0.7}/GaN$  MOS-HEMT at various bias conditions. The biasing operations at  $V_{gs} = -2$  V and  $-3$  V can be classified as class A and AB, respectively. At a bias of  $V_{gs} =$

$-3$  V and  $V_{ds} = 10$  V (class AB), a linear gain of 11.9 dB, a maximum output power of 20.26 dBm (531.44 mW/mm) and a maximum PAE of 26.43 %, were obtained. At higher  $V_{gs}$  ( $V_{gs} = -2$  V) and with  $V_{ds} = 10$  V (class A), a higher linear gain of 13.11 dB, a higher maximum output power of 22.3 dBm (849.73 mW/mm) and a lower maximum PAE of 23.77 % were achieved. At  $V_{gs} = -3$  V, the maximum output power increased from 531 mW/mm to 754 mW/mm with increased  $V_{ds}$  (from 10V to 15V), which is the same as for  $V_{gs} = -2$  V. These results show the potential for  $TiO_2/Al_{0.3}Ga_{0.7}/GaN$  MOS-HEMT to produce millimeter wavelength power.

Table 3: Power characteristics under various bias conditions.

$V_{gs}$	$V_{ds}$	$P_{out}$ Density (mW/mm)	Max PEA (%)	Linear gain (dB)
-2V	10V	577	23.6	12.7
-2V	15V	850	23.8	13.1
-3V	10V	531	26.5	11.9
-3V	15V	754	26.4	12.2

Table.4 demonstrates the performance of the proposed device by successful comparison of its effective channel thickness, peak  $f_r/f_{max}$ , and breakdown voltage with recent reported GaN MOS-HEMTs and HEMTs. In fact, by using the Regrown Source/Drain, the device exhibits a better  $I_{ds}$  at  $V_{gs} = 4$  V, an enhanced  $g_m$ , as well as higher  $f_r / f_{max}$  and breakdown voltage while compared to other MOS-HEMTs devices.

Table 4: Performance Comparison with other dielectric MOS-HEMTs

	This work	Ref[38]	Ref[23]	Ref[16]	Ref[39]	Ref[25]	Ref[40]	Ref[41]	Ref[13]	Ref[42]
Oxide type	$TiO_2$	---	$Al_2O_3$	$SiO_2$	$TiO_2$	$TiO_2$	$Al_2O_3$	$HfO_2$	$La_2O_3$	$TiO_2$
Gate length(nm)	60	40	210	100	2000	1000	1000	100	1000	50
Regrown Source/Drain	Yes	Yes	Yes	Yes	No	No	---	No	No	No
$f_r$ (GHz)	229	220	39.6	115	--	--	16.2	67	6.3	10.44
$f_{max}$ (GHz)	627	289	39.8	127	--	--	18.4	94	10.7	13.8
breakdown voltage(V)	477	42	114	--	490	139	126	--	131	132
$I_{ds}$ (mA)	632	1620	860	950	380	384	552	580	572	550
$g_m$ (mS/mm)	200	723	509	216	130	107	136	468	59.4	102
PAE (%)	26.43	--	--	--	--	--	16.4	--	--	--
sim/expr <sup>(i)</sup>	sim	expr	expr	expr	sim	expr	expr	sim	expr	expr

<sup>(i)</sup> sim/expr: simulation/experimental



al.

#### 4. Conclusion

The objective of this paper was to design and simulate a new MOS-HEMT model of an AlGa<sub>N</sub>/Ga<sub>N</sub> HEMT with 60 nm gate-length, high-k TiO<sub>2</sub> gate dielectric and Regrown Source/Drain. The obtained simulated results are very encouraging while compared to existing AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMTs (Table 4). This is, in fact, the best TiO<sub>2</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> MOS-HEMT simulated high-frequency performance reported so far, and thus, suitable for high-power RF circuit applications.

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